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Nizar Demni

Distributions of truncations of the heat kernel on the complex projective space

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Abstract

Let $(U_t)_{t \geq 0}$ be a Brownian motion valued in the complex projective space $\mathbb{C}P^{N-1}$. Using unitary spherical harmonics of homogeneous degree zero, we derive the densities of $|U_t^1|^2$ and of $(|U_t^1|^2, |U_t^2|^2)$, and express them through Jacobi polynomials in the simplices of $\mathbb{R}$ and $\mathbb{R}^2$ respectively. More generally, the distribution of $(|U_t^1|^2, \ldots, |U_t^k|^2)$, $2 \leq k \leq N-1$ may be derived using the decomposition of the unitary spherical harmonics under the action of the unitary group $\mathcal{U}(N-k+1)$ yet computations become tedious. We also revisit the approach initiated in [13] and based on a partial differential equation (hereafter pde) satisfied by the Laplace transform of the density. When $k = 1$, we invert the Laplace transform and retrieve the expression already derived using spherical harmonics. For general $1 \leq k \leq N-2$, integrations by parts performed on the pde lead to a heat equation in the simplex of $\mathbb{R}^k$.

1. Motivation

The complex unit sphere

$$S^{2N-1} = \{(z_1, \ldots, z_N), |z_1|^2 + \cdots + |z_N|^2 = 1\}, \quad N \geq 1,$$

is a compact manifold without boundary and therefore carries a Brownian motion $(U_t)_{t \geq 0}$ defined by means of its Laplace-Beltrami operator. This process is stationary and the random variable U_t converges weakly as $t \to \infty$ to a uniformly-distributed random vector U_∞. For the latter, it is already known that

$$(|U_\infty^1|^2, \ldots, |U_\infty^k|^2), \quad 1 \leq k \leq N-1,$$

follows the Dirichlet distribution ([9])

$$s_k(u)du = (1 - u_1 - u_2 - \cdots - u_k)^{N-k-1} \mathbf{1}_{\Sigma_k}(u) \prod_{i=1}^{k} du_i, \qquad (1.1)$$

where $\Sigma_k = \{u_i > 0, \ 1 \leq i \leq k, \ u_1 + \cdots + u_k < 1\}$ is the standard simplex. Motivated by quantum information theory, the investigations of the distribution of

$$U_t^{(k)} = (|U_t^1|^2, \ldots, |U_t^k|^2), \quad 1 \leq k \leq N,$$

started in [13] yet have not been completed. There, a linear pde for the Laplace transform of this distribution was obtained and partially solved only when $k = 1$. Recall that for the Brownian motion on the Euclidian sphere S^{N-1}, the density of a single coordinate is given by a series involving products of ultraspherical polynomials of index $(N-2)/2$ ([10]). The main ingredients leading to this series are the expansion of the heat kernel on S^{N-1} in the basis of $O(N)$-spherical harmonics and on Gegenbauer addition Theorem ([15], p.369). In the complex setting, it is very likely known that $(|U_t^1|^2)_{t \geq 0}$ is a real Jacobi process (see [3] and references therein). Nonetheless, one wonders how does the proof written in [10] carry to the Brownian motion on S^{2N-1} and how does it extend in order to derive the density of $U_t^{(k)}$. In the first part of this paper, we answer these questions by considering the heat kernel on the complex projective space $\mathbb{C}P^{N-1} = S^{2N-1}/S^1$ rather than S^{2N-1}. This is by no means a loss of generality since we are interested in the joint distribution of the moduli of k coordinates of U_t. Besides, the space of continuous functions on $\mathbb{C}P^{N-1}$ decomposes as the direct sum of subspaces of $\mathcal{U}(N)$-spherical harmonics that are homogeneous of degree zero, while the decomposition of continuous functions on S^{2N-1} involves all spherical harmonics ([8]). Accordingly, the heat kernel on $\mathbb{C}P^{N-1}$ is expressed as a series of normalized Jacobi polynomials $(P_n^{N-2,0}/P_n^{N-2,0}(1))_{n \geq 0}$ which for each n, gives the n-th reproducing kernel on $\mathbb{C}P^{N-1}$ ([11]). Hence, the integration over the sphere S^{2N-3} together with an application of Koornwinder's addition Theorem ([11]) lead to first result proved here:

Proposition 1.1. *The density of $|U_t^1|^2$ is given by*

$$f_t(c, u) = \left[\sum_{n=0}^{\infty} e^{-n(n+N-1)t} \frac{P_n^{N-2,0}(2c-1)P_n^{N-2,0}(2u-1)}{||P_n^{N-2,0}||_2^2} \right] s_1(u) \quad (1.2)$$

where we set $c = |U_0^1|^2 \in [0,1]$, and $||P_n||_2^2$ is the squared L^2-norm of $u \mapsto P_n^{N-2,0}(2u-1)$ with respect to $s_1(u)du$.

Up to an additional ingredient, the derivation of the density of $U_t^{(2)}$ is quite similar. Loosely speaking, we would like to integrate the heat kernel

over the sphere S^{2N-5} (we assume N large enough) and as such, we need to decompose degree zero homogeneous spherical harmonics in S^{2N-1} under the action of the unitary group $\mathcal{U}(N-1)$. This decomposition is stated in [12], Theorem 5.1, and the n-th reproducing kernel in turn decomposes as a weighted sum of reproducing kernels on S^{2N-3}. Consequently, Koornwinder's addition Theorem again leads to the sought density which may be expressed through Jacobi polynomials in the simplex Σ_2 ([6], Proposition 2.3.8 p.47). More precisely

Proposition 1.2. *Let* $(Q_{j,n-j}^{(N)})_{n \geq 0, 0 \leq j \leq n}$ *denote the family of Jacobi polynomials in the simplex* Σ_2. *Then the density of* $U_t^{(2)}$ *reads*

$$\left[\sum_{n=0}^{\infty} e^{-n(n+N-1)t} \sum_{j=0}^{n} \frac{Q_{n-j,j}^{(N)}(c_1, c_2) Q_{n-j,j}^{(N)}(u_1, u_2)}{||Q_{n-j,j}^{(N)}||_2^2} \right] s_2(u_1, u_2), \qquad (1.3)$$

where we set $(c_1, c_2) = (|U_0^1|^2, |U_0^2|^2) \in \Sigma_2$ *and* $||Q_{n-j,j}^{(N)}||_2^2$ *is the squared* L^2-*norm of* $Q_{n-j,j}^{(N)}$ *with respect to* $s_2(u_1, u_2)du_1du_2$.

More generally, the derivation of the distribution of $U_t^{(k)}, 2 \leq k \leq N-1$ relies on the decomposition of the spherical harmonics under the action of $\mathcal{U}(N-k+1)$. The resulting density with respect to Lebesgue measure du is expressed through orthonormal Jacobi polynomials in Σ_k as

$$\sum_{n \geq 0} e^{-n(n+N-1)t} \sum_{\tau \in \mathbb{N}^k, |\tau|=n} Q_\tau^{(N)}(c_1, c_2, \ldots, c_k) Q_\tau^{(N)}(u) s_k(u)$$

where $(c_1, \ldots, c_k) = (|U_0^1|^2, \ldots, |U_0^k|^2)$. Yet computations become tedious and we are not willing to exhibit them here. Rather, we shall revisit and complete the investigations started in [13]. Actually, an expression for the Laplace transform of the density of $|U_t^1|^2$ was obtained there and involves the following sequence $(a_n = a_n(c, N))_{n \geq 0}$ of real numbers determined recursively by ([13], eq. 4.23)

$$\sum_{n=0}^{p} a_n \binom{p}{n} \frac{1}{(N+2n)_{p-n}} = \frac{c^p}{p!}, \quad p \geq 0, \qquad (1.4)$$

where $a_0 = 1$ and $(x)_p = \Gamma(x+p)/\Gamma(p)$ is the Pochhammer symbol. In particular, the following was proved ([13] eq. 4.24. and eq. 4.25):

$$a_n(0, N) = \frac{(-1)^n}{(N+n-1)_n}, \quad a_n(1, N) = \frac{(N-1)_n}{n!(N+n-1)_n},$$

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which we can rewrite as

$$\frac{1}{(N+n-1)_n}P_n^{N-2,0}(-1), \quad \frac{1}{(N+n-1)_n}P_n^{N-2,0}(1)$$

respectively. Using a Neumann series for Bessel functions ([15]), we shall prove that

Proposition 1.3. *For any* $c \in [0,1]$,

$$a_n = a_n(c,N) = \frac{1}{(N+n-1)_n}P_n^{N-2,0}(2c-1). \tag{1.5}$$

Having these coefficients in hands, we can then invert the Laplace transform and retrieve (1.2). At this level, we point out that the pde satisfied by the Laplace transform of the density of $|U_t^1|^2$ leads after integrations by parts to the heat equation associated with the Jacobi operator

$$u(1-u)\partial_u^2 + [1 - Nu]\partial_u. \tag{1.6}$$

More generally, the pde satisfied by the Laplace transform of the joint distribution of $U_t^{(k)}$, $1 \le k \le N-2$, gives rise to the heat equation on the standard simplex associated with the generalized Jacobi operator ([1], see also [6] p.46 but consult the list of errata available on the webpage of Y. Xu):

$$\sum_{i=1}^{k}[1 - Nu_i]\partial_i + \sum_{i=1}^{k}(u_i - u_i^2)\partial_{ii} - \sum_{i \ne j}u_i u_j \partial_{ij}. \tag{1.7}$$

This is an elliptic operator admitting different orthogonal basis of eigenpolynomials corresponding to the sequence of eigenvalues $\{-n(N+n-1), n \ge 0\}$. Among them figure the Jacobi polynomials in the simplex, which agrees with our previous computations.

The paper is organized as follows. The two following sections are devoted to the derivations of the densities displayed in (1.2) and (1.3). In section 4, we prove proposition 1.3 and invert the Laplace transform of the density of $|U_t^1|^2$. In section 5, we perform integrations by parts on the pde satisfied by the Laplace transform of the density of $U_t^{(k)}$, omitting for a while the boundary terms. In the last section, we write down the latters and show that all of them vanish unless $k = N-1$.

Remark 1.4. The fact that $U_t^{(k)}$ coincides with the random variable studied in [13] is justified as follows. The sphere $S^{2N-1} \approx \mathcal{U}_N/\mathcal{U}_{N-1}$ is a homogeneous space and if $\rho : \mathcal{U}_N \to \mathcal{U}_N/\mathcal{U}_{N-1}$ is the quotient map then Lemma

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2.1. in [4] implies that

$$\Delta_{\mathcal{U}_N}(f \circ \rho) = \Delta_{\mathcal{U}_N/\mathcal{U}_{N-1}}(f) \circ \rho$$

holds for any smooth function f on $\mathcal{U}_N/\mathcal{U}_{N-1}$.
Similarly, if $h : S^{2N-1} \to \mathbb{C}P^{N-1}$ is the Hopf projection then we also have from Proposition G. III. 15 in [5]:

$$\Delta_{S^{2N-1}}(f \circ h) = \Delta_{\mathbb{C}P^{N-1}}(f) \circ h$$

for any smooth function f on $\mathbb{C}P^{N-1}$.

2. **The heat kernel on $\mathbb{C}P^{N-1}$ and the distribution of $|U_t^1|^2$**

Let m, n be non negative integers and recall from [11] that (m, n)-complex spherical harmonics are the restriction to S^{2N-1} of harmonic polynomials in the variables

$$(z_1, z_2, \ldots, z_N, \overline{z_1}, \overline{z_2}, \ldots, \overline{z_N})$$

which are m-homogenous in the variables $(z_i)_{i=1}^N$ and n-homogeneous in the variables $(\overline{z_i})_{i=1}^N$. Taking $m = n$, we obtain the (n, n)-complex spherical harmonics that are homogenous of degree zero with respect to the action of S^1. Their restrictions to $\mathbb{C}P^{N-1}$ form a dense algebra in the space of continuous functions on $\mathbb{C}P^{N-1}$ endowed with the uniform norm (see [8], p.189). Moreover, the spectrum of the Laplace-Beltrami operator on $\mathbb{C}P^{N-1}$ is given by the sequence $\{-n(n + N - 1), n \geq 0\}$[1]. Hence, the corresponding heat kernel is expanded in any orthonormal (with respect to the volume measure $\mathrm{vol}_{\mathbb{C}P^{N-1}}$) basis of homogeneous spherical harmonics of degree zero $(Y_j)_{j \geq 1}$ as:

$$R_t(w, z) = \sum_{n=0}^{\infty} e^{-n(n+N-1)t} \sum_{j=1}^{d(n,N)} Y_j(w)\overline{Y_j}(z), \quad w, z \in \mathbb{C}P^{N-1}.$$

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[1]We normalize the Laplacian on $\mathbb{C}P^{N-1}$ by a factor $1/4$.

Here $d(n, N)$ is the dimension of the eigenspace of (n, n)-complex spherical harmonics given by (Theorem 3.6 in [11])

$$d(n, N) = \frac{2n + N - 1}{N - 1} \left(\frac{(N-1)_n}{n!} \right)^2.$$

Besides, the reproducing kernel formula (Theorem 3.8 in [11][2]) shows that the kernel R_t is real and does not depend on the choice of the basis (this is the analogue of (22) in [10]):

$$R_t(w, z) = 2\pi \sum_{n=0}^{\infty} e^{-n(n+N-1)t} \frac{d(n, N)}{\text{vol}(S^{2N-1})} \frac{P_n^{N-2,0}(2|\langle w, z \rangle|^2 - 1)}{P_n^{N-2,0}(1)}, \quad (2.1)$$

where $P_n^{N-2,0}$ is the n-th Jacobi polynomial of parameters $(N-2, 0)$ ([2], p.295),

$$P_n^{N-2,0}(1) = \frac{(N-1)_n}{n!}, \quad \langle w, z \rangle = \sum_{i=1}^{N} w_i \overline{z_i},$$

and

$$\text{vol}(S^{2N-1}) = \frac{2\pi^N}{(N-1)!}.$$

is the volume of S^{2N-1}. Note that

$$R_t(w, z) = \frac{(N-2)!}{\pi^{N-1}} \sum_{n=0}^{\infty} e^{-n(n+N-1)t} \frac{P_n^{N-2,0}(1) P_n^{N-2,0}(2|\langle w, z \rangle|^2 - 1)}{||P_n^{N-2,0}||_2^2}$$

where

$$||P_n^{N-2,0}||_2^2 = \frac{1}{2n + N - 1}$$

is the squared L^2-norm of $u \mapsto P_n^{N-2,0}(2u - 1)$ with respect to $s_1(u)du$ ([2], p.99). Thus Gasper's Theorem entails the positivity of R_t ([7]). Now, we proceed to the derivation of the density of $|U_t^1|^2$ and start with the decompositions

$$
\begin{aligned}
w &= \cos\theta_1 e_1 + \sin\theta_1 \xi_1, \\
z &= \cos\theta_2 e_1 + \sin\theta_2 \xi_2
\end{aligned}
$$

where e_1 is the first vector of the canonical basis of $\mathbb{C}^N$, $\theta_1, \theta_2 \in (0, \pi/2)$, $\xi_1, \xi_2 \in S^{2N-3}$. The volume measure of $\mathbb{C}P^{N-1}$ in turn splits as (see [11],

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[2]The additional factor 2π comes from the fact that $\text{vol}(S^{2N-1}) = 2\pi\text{vol}(\mathbb{C}P^{N-1})$.

eq.2.18)

$$\text{vol}_{\mathbb{C}P^{N-1}}(dz) = \cos\theta_2(\sin\theta_2)^{2N-3}d\theta_2\text{vol}_{S^{2N-3}}(d\xi_2).$$

and the next step is to integrate (2.1) over ξ_2. But $U(N-1)$ acts transitively on S^{2N-3} therefore we can take $\xi_1 = e_2$ to be the second vector of the canonical basis. As such, we are left with the volume of S^{2N-5} (if N is large enough) and with the integration over the distribution of the first coordinate of ξ_2. If this coordinate is parametrized by (r, ψ) then its distribution reads

$$r(1-r^2)^{N-3}\mathbf{1}_{[0,1]}(r)\mathbf{1}_{[0,2\pi]}(\psi)dr\,d\psi.$$

Consequently, the density of $|U_t^1|^2$ displayed in (1.2) follows from the product formula (4.12) in [11] together with the variables change $u = \cos^2\theta_2$ $(c = \cos^2\theta_1)$.

Remark 2.1. The eigenvalue of a (n, n)-spherical harmonic equals the eigenvalue of a $O(2N)$-spherical harmonic of degree $2n$ in S^{2N-1} viewed as a real Euclidian sphere. This coincidence is due to the fact that both polynomials are homogenous with the same total degree $2n$ and since the corresponding eigenvalue comes from the action of the Euler operator

$$\sum_{i=1}^{N} z_i\partial_{z_i} + \sum_{i=1}^{N} \overline{z_i}\partial_{\overline{z_i}} = \sum_{i=1}^{N} x_i\partial_{x_i} + \sum_{i=1}^{N} y_i\partial_{y_i},$$

where $z_i \in \mathbb{C}$ is identified with $(x_i, y_i) \in \mathbb{R}^2$.

3. **The distribution of** $(|U_t^1|^2, |U_t^2|^2)$

Up to an additional ingredient, the lines of the previous proof enable to derive the density of $U_t^{(2)}$. More precisely, we start with the decompositions

$$
\begin{aligned}
w &= \cos\theta_1 e_1 + \sin\theta_1\xi_1 \\
&= \cos\theta_1 e_1 + \sin\theta_1\cos\beta_1 e^{i\phi_1}e_2 + \sin\theta_1\sin\beta_1\eta_1, \\
z &= \cos\theta_2 e_1 + \sin\theta_2\xi_2 \\
&= \cos\theta_2 e_1 + \sin\theta_2\cos\beta_2 e^{i\phi_2}e_2 + \sin\theta_2\sin\beta_2\eta_2,
\end{aligned}
$$

where $\beta_1, \beta_2 \in (0, \pi/2), \phi_1, \phi_2 \in (0, 2\pi), \eta_1, \eta_2 \in S^{2N-5}$ and e_2 is the second vector of the canonical basis of $\mathbb{C}^N$. We also split the volume measure

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on $\mathbb{C}P^{N-1}$ as

$$\mathrm{vol}_{\mathbb{C}P^{N-1}}(dz) =$$
$$\left(\cos\theta_2\sin^{2N-3}(\theta_2)\cos\beta_2\sin^{2N-5}(\beta_2)d\theta_2d\beta_2d\phi_2\right)\mathrm{vol}_{S^{2N-5}}(d\eta_2).$$

Now comes the needed additional ingredient, which is the special instance $m = n, \phi_1 = \phi_2 = 0$ in the formula stated in the bottom of p.5 in [12]. In order to recall it, let

$$p_j^{a,b}(x) = \frac{P_j^{a,b}(x)}{P_j^{a,b}(1)}, \quad a,b > -1,$$

be the j-th normalized Jacobi polynomial and define the complex-valued polynomial ([11], eq.3.15)

$$R_{j,q}^{\alpha}(y) = |y|^{|j-q|}e^{i(j-q)\arg(y)}p_{j\wedge q}^{\alpha,|j-q|}(2|y|^2 - 1), \quad y \in \mathbb{C}, \alpha > -1,$$

as well as ([12], p.6)

$$c_{j,q}(n,N) = \frac{N-2}{N-2+q+j}\binom{n}{q}\binom{n}{j}\frac{(N+n-1)_q(N+n-1)_j}{(N-2+j)_q(N-2+q)_j}.$$

Then the n-th reproducing kernel on $\mathbb{C}P^{N-1}$ admits the following expansion

$$\frac{P_n^{N-2,0}(2|\langle w,z\rangle|^2 - 1)}{P_n^{N-2,0}(1)} =$$
$$\sum_{j,q=0}^{n} c_{j,q}(n,N)[\sin\theta_1\sin\theta_2]^{j+q}[\cos\theta_1\cos\theta_2]^{|j-q|}p_{n-j\wedge n-q}^{N-2+j+q,|j-q|}(\cos 2\theta_1)$$
$$p_{n-j\wedge n-q}^{N-2+j+q,|j-q|}(\cos 2\theta_2)R_{j,q}^{N-3}\left(\langle\xi_1,\xi_2\rangle\right).$$

Substituting in (2.1), we see that the next step towards the joint distribution of (U_t^1, U_t^2) consists in integrating

$$R_{j,q}^{N-3}\left(\langle\xi_1,\xi_2\rangle\right) = R_{j,q}^{N-3}\left(\cos\beta_1\cos\beta_2 + \sin\beta_1\sin\beta_2\langle\eta_1,\eta_2\rangle\right)$$

over $\eta_2 \in S^{2N-3}$. To this end, we can assume without loss of generality that $\eta_1 = e_3$ (the third vector of the canonical basis) and use formula

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(4.11) in [11]. Altogether, we get

$$\frac{2\pi^2\mathrm{vol}(S^{2N-5})}{(N-3)\mathrm{vol}(S^{2N-1})}\sum_{n=0}^{\infty}e^{-n(n+N-1)t}d(n,N)\sum_{j,q=0}^{n}c_{j,q}(n,N)[\sin\theta_1\sin\theta_2]^{j+q}[\cos\theta_1\cos\theta_2]^{|j-q|}$$

$$p_{n-j\wedge n-q}^{N-2+j+q,|j-q|}(\cos 2\theta_1)p_{n-j\wedge n-q}^{N-2+j+q,|j-q|}(\cos 2\theta_2)R_{j,q}^{N-3}(\cos\beta_1 e^{i\phi_1})R_{j,q}^{N-3}(\cos\beta_2 e^{i\phi_2})$$

$$=\frac{2(N-2)}{\pi}\sum_{n=0}^{\infty}e^{-n(n+N-1)t}(2n+N-1)[P_n^{N-2,0}(1)]^2\sum_{j,q=0}^{n}c_{j,q}(n,N)[\sin\theta_1\sin\theta_2]^{j+q}$$

$$[\cos\theta_1\cos\theta_2]^{|j-q|}\,p_{n-j\wedge n-q}^{N-2+j+q,|j-q|}(\cos 2\theta_1)\,p_{n-j\wedge n-q}^{N-2+j+q,|j-q|}(\cos 2\theta_2)$$

$$p_{j\wedge q}^{N-3,|j-q|}(\cos 2\beta_1)p_{j\wedge q}^{N-3,|j-q|}(\cos 2\beta_2)e^{i(j-q)(\phi_1+\phi_2)}$$

with respect to

$$\cos\theta_2\sin^{2N-3}\theta_2\cos\beta_2\sin^{2N-5}\beta_2 d\theta_2 d\beta_2 d\phi_2.$$

Integrating over $\phi_2 \in (0, 2\pi)$, then the sum over (j,q) reduces to a sum over $q = j$. Thus, the density of (θ_2, β_2) given (θ_1, β_1) reads

$$4(N-2)\sum_{n=0}^{\infty}e^{-n(n+N-1)t}(2n+N-1)[P_n^{N-2,0}(1)]^2\sum_{j=0}^{n}c_{j,j}(n,N)[\sin\theta_1\sin\theta_2]^{2j}$$

$$p_{n-j}^{N-2+2j,0}(\cos 2\theta_1)p_{n-j}^{N-2+2j,0}(\cos 2\theta_2)p_j^{N-3,0}(\cos 2\beta_1)p_j^{N-3,0}(\cos 2\beta_2)$$

with respect to

$$\cos\theta_2\sin^{2N-3}\theta_2\cos\beta_2\sin^{2N-5}\beta_2 d\theta_2 d\beta_2.$$

Performing the variables change

$$u=\cos\theta_2, v=\sin\theta_2\cos\beta_2,$$

we deduce that the density of $(|U_t^1|, |U_t^2|)$ given $(|U_0^1|, |U_0^2|)$ is:

$$4(N-2)\sum_{n=0}^{\infty}e^{-n(n+N-1)t}(2n+N-1)[P_n^{N-2,0}(1)]^2\sum_{j=0}^{n}c_{j,j}(n,N)[(1-|U_0^1|^2)(1-u^2)]^j$$

$$p_{n-j}^{N-2+2j,0}(2|U_0^1|^2-1)\,p_{n-j}^{N-2+2j,0}(2u^2-1)\,p_j^{N-3,0}\left(\frac{2|U_0^2|^2}{1-|U_0^1|^2}-1\right)p_j^{N-3,0}\left(\frac{2v^2}{1-u^2}-1\right)$$

with respect to

$$uv(1-u^2-v^2)^{N-3}\mathbf{1}_{\{u>0,v>0,u^2+v^2<1\}}du\,dv.$$

Finally, if $(|U_0^1|^2, |U_0^2|^2) = (c_1, c_2) \in \Sigma_2$ then the density of $U_t^{(2)}$ reads

$$(N-2)\sum_{n=0}^{\infty} e^{-n(n+N-1)t}(2n+N-1)[P_n^{N-2,0}(1)]^2 \sum_{j=0}^{n} c_{j,j}(n,N)[(1-c_1)(1-u_1)]^j$$

$$p_{n-j}^{N-2+2j,0}(2c_1-1)\, p_{n-j}^{N-2+2j,0}(2u_1-1)\, p_j^{N-3,0}\left(\frac{2c_2}{1-2c_1}-1\right) p_j^{N-3,0}\left(\frac{2u_2}{1-u_1}-1\right)$$

with respect to $s_2(u_1, u_2)du_1 du_2$. In order to get (1.3), set

$$Q_{n-j,j}^{(N)}(u_1, u_2) = (1-u_1)^j P_{n-j}^{N-2+2j,0}(2u_1-1)P_j^{N-3,0}\left(\frac{2u_2}{1-u_1}-1\right),$$
$$n \geq 0, j = 0, 1, \ldots, n.$$

These are Jacobi polynomials in the simplex Σ_2 and are orthogonal with respect to the Dirichlet distribution whose density is $s_2(u_1, u_2)$ (specialize Proposition 2.3.8 in [6] to $\alpha = (n-j, j), \kappa = (1/2, 1/2, N-5/2)^3$). After some computations, the density of $U_t^{(2)}$ may be written as

$$\sum_{n=0}^{\infty} e^{-n(n+N-1)t} \sum_{j=0}^{n} \frac{Q_{n-j,j}^{(N)}(c_1,c_2)Q_{n-j,j}^{(N)}(u_1,u_2)}{\|Q_{n-j,j}^{(N)}\|_2^2}$$

where

$$\begin{aligned}\|Q_{n-j,j}^{(N)}\|_2^2 &= \frac{1}{(2n+N-1)(2j+N-2)} \\ &= \frac{[P_{n-j}^{N-2+2j,0}(1)P_j^{N-3,0}(1)]^2}{(N-2)(2n+N-1)[P_n^{N-2,0}(1)]^2 c_{j,j}(n,N)}\end{aligned}$$

is the squared L^2-norm of $Q_{n-j,j}^{(N)}$ with respect to $s_2(u_1,u_2)du_1 du_2$.

Remark 3.1. For general $k \geq 3$, the density of $U_t^{(k)}$ may be derived in a similar way by decomposing the variable $z \in \mathbb{C}P^{N-1}$ and the spherical harmonics on S^{2N-1} over the sphere $S^{2N-2k+1}$. From the point of view of representation theory, this is equivalent to the decompositions of the representation of $\mathcal{U}(N)$ in the space of $\mathcal{U}(N)$-spherical harmonics under the action of the subgroup $\mathcal{U}(N-k+1)$. A similar statement is valid when one considers $O(N)$-spherical harmonics (see [14], Ch.IX).

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[3] Beware of the different normalization of the Jacobi polynomials used in Proposition 2.3.8. The reader is also invited to consult the list of errata available on the webpage of Y. Xu.

4. **The distribution of $|U_t^1|^2$: another proof**

In this section, we shall prove proposition 1.3. To this end, we rewrite (1.4) as

$$\sum_{n=0}^{p} a_n \frac{1}{n!(p-n)!} \frac{1}{(N+2n)_{p-n}} = \frac{c^p}{(p!)^2}, \qquad (4.1)$$

multiply both sides of (4.1) by $(-1)^p(x/2)^{2p+N-1}$ for x lying in some neighborhood of zero then sum over $p \geq 0$. Interchanging the order of summation, the system (1.4) is equivalent to

$$\sum_{n \geq 0} \frac{a_n}{n!} \Gamma(N+2n)(-1)^n J_{2n+N-1}(x) = J_0(\sqrt{c}x) \left(\frac{x}{2}\right)^{N-1}$$

where J_α is the Bessel function of index $\alpha \in \mathbb{R}$ defined by ([15]):

$$J_\alpha(x) = \sum_{p \geq 0} \frac{(-1)^p}{p!\,\Gamma(p+\alpha+1)} \left(\frac{x}{2}\right)^{2p+\alpha}.$$

Note in passing that the estimate ([15])

$$|\Gamma(N+2n)J_{2n+N-1}(x)| \leq \left(\frac{|x|}{2}\right)^{2n+N-1}$$

shows that (4.1) converges provided

$$\sum_{n \geq 0} \frac{|a_n|}{n!} \left(\frac{|x|}{2}\right)^{2n} \qquad (4.2)$$

does. Now recall the Neumann series ([15], p.138)

$$\left(\frac{x}{2}\right)^\nu = \sum_{n \geq 0} \frac{(\nu+2n)\Gamma(\nu+n)}{n!} J_{\nu+2n}(x), \quad \nu \in \mathbb{N}.$$

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Specializing it to $\nu = N - 2$, we get

$$(x/2)^{N-1} J_0(\sqrt{c}x)$$

$$= \sum_{p \geq 0} \frac{(-1)^p c^p}{(p!)^2} \sum_{n \geq 0} \frac{(2p+2n+N-1)\Gamma(2p+n+N-1)}{n!} J_{2p+2n+N-1}(x)$$

$$= \sum_{p \geq 0} \frac{(-1)^p c^p}{(p!)^2} \sum_{n \geq p} \frac{(2n+N-1)\Gamma(p+n+N-1)}{(n-p)!} J_{2n+N-1}(x)$$

$$= \sum_{n \geq 0} (2n+N-1) \sum_{p=0}^{n} \frac{(-1)^p c^p}{(p!)^2} \frac{\Gamma(p+n+N-1)}{(n-p)!} J_{2n+N-1}(x).$$

Substituting in (4.1), then the uniqueness of the solution of (1.4) yields

$$\frac{a_n}{n!} \Gamma(N+2n)(-1)^n = (2n+N-1) \sum_{p=0}^{n} \frac{c^p}{(p!)^2} \frac{\Gamma(p+n+N-1)}{(n-p)!}$$

or equivalently

$$a_n = \frac{(-1)^n n!}{\Gamma(N+2n-1)} \sum_{p=0}^{n} \frac{(-1)^p c^p}{(p!)^2} \frac{\Gamma(p+n+N-1)}{(n-p)!}$$

$$= \frac{(-1)^n}{(N+n-1)_n} \sum_{p=0}^{n} \frac{(-1)^p n!}{(n-p)!} \frac{(n+N-1)_p}{(p!)^2} c^p$$

$$= \frac{(-1)^n}{(N+n-1)_n} {}_2F_1(-n, n+N-1, 1, c)$$

where ${}_2F_1$ is the Gauss hypergeometric function ([2]). But from the very definition of Jacobi polynomials ([2], p.99)

$$P_n^{\alpha,\beta}(x) = \frac{(\alpha+1)_n}{n!} {}_2F_1(-n, n+\alpha+\beta+1, \alpha+1, (1-x)/2), \ \alpha, \beta > -1,$$

and the relation $P_n^{\alpha,\beta}(x) = (-1)^n P_n^{\beta,\alpha}(-x)$ ([2], p.305), we obtain (1.5) as required.

Remark 4.1. The estimate

$$|P_n^{0,N-2}(1-2c)| \leq |P_n^{0,N-2}(-1)|$$

which follows for instance from the integral representation of the Gauss hypergeometric function shows that the series (4.2) indeed converges absolutely everywhere.

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With (1.5) in hands, we can invert the Laplace transform ([13], eq.4.11):

$$\varphi_{t/N}(c, \lambda) = \int_0^1 e^{\lambda u} f_{t/N}(c, u) du = \sum_{n=0}^{\infty} a_n e^{-\Lambda_n t} \lambda^n {}_1F_1(n+1, N+2n, \lambda), \quad \lambda \in \mathbb{R},$$

where ${}_1F_1$ is the confluent hypergeometric function ([2], [15]) and $\Lambda_n = n(n + N - 1)/N$. To proceed, recall the integral representation ([2], p.234)

$$ {}_1F_1(a, b, \lambda) = \frac{\Gamma(b)}{\Gamma(b-a)\Gamma(a)} \int_0^1 e^{\lambda u} u^{a-1} (1-u)^{b-a-1} du, \quad b > a > 0.$$

It follows that

$$\lambda^n {}_1F_1(n+1, N+2n, \lambda) = \frac{\Gamma(N+2n)}{\Gamma(N+n-1)n!} \int_0^1 \lambda^n e^{\lambda u} u^n (1-u)^{N+n-2} du$$

$$= \frac{(-1)^n \Gamma(N+2n)}{\Gamma(N+n-1)n!} \int_0^1 e^{\lambda u} \left(\frac{d}{du}\right)^n [u^n (1-u)^{N+n-2}] du$$

after n integration by parts. But Rodriguez formula ([2], p.99)

$$(1-x)^{\alpha}(1+x)^{\beta} P_n^{\alpha,\beta}(x) = \frac{(-1)^n}{2^n n!} \left(\frac{d}{du}\right)^n [(1-x)^{n+\alpha}(1+x)^{n+\beta}], \quad x \in (-1, 1)$$

together with the variable change $x = 1 - 2u$ yields

$$\left(\frac{d}{du}\right)^n [u^n (1-u)^{N+n-2}] = n!(1-u)^{N-2} P_n^{0,N-2}(1-2u).$$

As a result

$$a_n \lambda^n {}_1F_1(n+1, N+2n, \lambda) = P_n^{0,N-2}(1-2c) \int_0^1 e^{\lambda u} P_n^{0,N-2}(1-2u)(1-u)^{N-2} du$$

and Tonelli-Fubini Theorem yields (1.2) at time t/N.

5. From the Laplace transform to the generalized Jacobi operator

Another way to come from $\varphi_t(c, \lambda)$ to $f_t(c, \lambda)$ is as follows. For sake of simplicity, we shall drop the dependence on the parameter c. So, recall from [13] Proposition 4.2 that φ satisfies

$$\partial_t \varphi = \lambda \varphi + (\lambda^2 - N\lambda) \partial_\lambda \varphi - \lambda^2 \partial_\lambda^2 \varphi$$

with the initial conditions $\varphi_0(c,\lambda) = e^{\lambda c}, \varphi_t(c,0) = 1$. Assume the density f is unknown and is smooth in both variables (t,u), then integration by parts yield

$$\int_0^1 e^{\lambda u} \partial_t f_t(u) du = [e^{\lambda u}(1-Nu)f_t(u)]_0^1 + \int_0^1 e^{\lambda u} \partial_u[(Nu-1)f_t(u)]du$$

$$+ [\lambda e^{\lambda u}u(1-u)f_t(u)]_0^1 - [e^{\lambda u}\partial_u(u(1-u)f_t(u))]_0^1$$

$$+ \int_0^1 e^{\lambda u} \partial_u^2[u(1-u)f_t(u)]du$$

$$= e^{\lambda}(2-N)f_t(1) + \int_0^1 e^{\lambda u} \mathscr{L}(f_t)(u)du$$

where

$$\mathscr{L} = u(1-u)\partial_u^2 + [1 + (N-4)u]\partial_u + (N-2).$$

If $N = 2$ then

$$\mathscr{L} = u(1-u)\partial_u^2 + [1-2u]\partial_u$$

is nothing else but (1.6) with $N = 2$. Otherwise, write $f_t(u) = g_t(u)s_1(u)$ for a smooth function g and note that $\mathscr{L}(s_1) = 0$. As a result,

$$\mathscr{L}(f_t)(u) = s_1(u)\left\{u(1-u)\partial_u^2 + [1 - Nu]\partial_u\right\}(g_t)(u)$$

where the RHS is the operator displayed in (1.6). Since $f_t(1) = 0$ when $N \geq 3$ then we always have

$$\int_0^1 e^{\lambda u} \partial_t g_t(u)s_1(u)du = \int_0^1 e^{\lambda u}\left\{u(1-u)\partial_u^2 + [1 - Nu]\partial_u\right\}(g_t)(u)s_1(u)du.$$

But the set of monomials $(u^n)_{n\geq 0}$ is total in $L^2([0,1], (1-u)^{N-2}du)$ ([6], Theorem 3.17) then g solves the heat equation

$$\partial_t g_t(u) = \left\{u(1-u)\partial_u^2 + [1 - Nu]\partial_u\right\}(g_t)(u).$$

More generally, the Laplace transform of the density of $U_t^{(k)}, 1 \leq k \leq N-1$ satisfies the linear pde

$$\partial_t \varphi \quad = \quad \sum_{j=1}^k \lambda_j \varphi + \sum_{j=1}^k (\lambda_j^2 - N\lambda_j)\partial_j \varphi - \sum_{j,i=1}^k \lambda_i \lambda_j \partial_{ij}\varphi. \qquad (5.1)$$

Hoping there will be no confusion, set again

$$\varphi_t(c,\lambda) = \int_{\Sigma} e^{\langle \lambda, u \rangle} f_t(c,u)du$$

where du is the Lebesgue measure in the simplex Σ. Then, integration by parts

$$\sum_{i=1}^{k} \lambda_i \varphi_t(\lambda) \quad \to \quad -\int_{\Sigma} e^{\langle \lambda, u \rangle} \left(\sum_{i=1}^{k} \partial_i f_t(u) \right) du$$

$$(\lambda_i^2 - N\lambda_i)\partial_i \varphi_t(\lambda) \quad \to \quad \int_{\Sigma} e^{\langle \lambda, u \rangle} [\partial_{ii} + N\partial_i](u_i f_t(u)) du, \quad i \in \{1, \ldots, k\},$$

$$\lambda_i^2 \partial_{ii} \varphi_t(\lambda) \quad \to \quad \int_{\Sigma} e^{\langle \lambda, u \rangle} \partial_{ii}(u_i^2 f_t(u)) du, \quad i \in \{1, \ldots, k\},$$

$$\lambda_i \lambda_j \partial_{ij} \varphi_t(\lambda) \quad \to \quad \int_{\Sigma} e^{\langle \lambda, u \rangle} \partial_{ij}(u_i u_j f_t(u)) du, \quad 1 \le i \ne j \le k,$$

transform the pde (5.1) into

$$\int_{\Sigma} e^{\langle \lambda, u \rangle} \partial_t f_t(u) du = \text{boundary terms} + \int_{\Sigma} e^{\langle \lambda, u \rangle} \mathscr{L}(f_t)(u) du,$$

where this time $\mathscr{L}$ denotes the operator

$$k(N - k - 1) + \sum_{i=1}^{k} [1 + [N - 4 - 2(k-1)]u_i]\partial_i + \sum_{i=1}^{k} (u_i - u_i^2)\partial_{ii} - \sum_{i \ne j} u_i u_j \partial_{ij}.$$

If $k = N - 1$ then $\mathscr{L}$ reduces to the operator displayed in (1.7). Otherwise, set

$$f_t(u) = g_t(u)s_k(u)$$

where this time g is a smooth function in both variables $(t, u), u \in \Sigma_k$ and note that the relations $\partial_i s_k = \partial_1 s_k, 1 \le i \le k$ together with the identity

$$\sum_{i=1}^{k} u_i(1 - u_i) - \sum_{i \ne j} u_i u_j = \sum_{i=1}^{k} u_i \left(1 - \sum_{i=1}^{k} u_i \right) = \left(1 - \sum_{i=1}^{k} u_i \right) - \left(1 - \sum_{i=1}^{k} u_i \right)^2$$

imply that $\mathscr{L}_u(s_k) = 0$. Hence $\mathscr{L}(f_t)(u)$ gives rise to l

$$\sum_{i=1}^{k} [1 + [N - 4 - 2(k-1)]u_i]\partial_i + \sum_{i=1}^{k}(u_i - u_i^2)\partial_{ii}$$

$$- \sum_{i \neq j} u_i u_j \partial_{ij} + 2\frac{\partial_1 s}{s}(u)\left\{\sum_{i=1}^{k} u_i(1 - u_i)\partial_i - \sum_{i \neq j} u_j u_i \partial_i\right\}$$

$$= \sum_{i=1}^{k} [1 + [N - 4 - 2(k-1)]u_i]\partial_i + \sum_{i=1}^{k}(u_i - u_i^2)\partial_{ii}$$

$$- \sum_{i \neq j} u_i u_j \partial_{ij} + 2\frac{\partial_1 s}{s}(u)\left\{\sum_{i=1}^{k} u_i\partial_i\left(1 - u_i - \sum_{j \neq i} u_j\right)\right\}$$

$$= \sum_{i=1}^{k} [1 - Nu_i]\partial_i + \sum_{i=1}^{k}(u_i - u_i^2)\partial_{ii} - \sum_{i \neq j} u_i u_j \partial_{ij}$$

acting on g_t. Consequently, if the boundary terms vanish then Theorem 3.17 in [6] implies that g_t solves the heat equation

$$\left[\sum_{i=1}^{k} [1 - Nu_i]\partial_i + \sum_{i=1}^{k}(u_i - u_i^2)\partial_{ii} - \sum_{i \neq j} u_i u_j \partial_{ij}\right] g_t = \partial_t g_t.$$

We shall see below that this is the case provided that $1 \leq k \leq N - 2$.

6. **Analysis of the boundary terms**

Recall from the previous section that the integration by parts performed in the one-variable setting gave rise to the boundary term

$$e^{\lambda u}\left[(1 - Nu)f_t(u) - \partial_u(u(1-u)f_t(u))\right]$$

which vanishes at $u = 1$ since $f_t(1) = 0$ when $N \geq 3$ (note that there is no such condition when $N = 2$). For higher values $k \geq 2$, the situation is similar provided that $1 \leq k \leq N - 2$ and is different when $k = N - 1$ due to the interactions

between u_i and u_j for $i \neq j$. Indeed, the boundary terms are given by

$$\sum_{i=1}^{k} \int [e^{\lambda_i u_i}(1-Nu_i)f_t(u)]_0^{1-\sum_{j\neq i} u_j} \left(\prod_{j\neq i} e^{\lambda_j u_j} \mathbf{1}_{[0,1]}(u_j)du_j \right) +$$

$$\sum_{i=1}^{k} \int [\lambda_i e^{\lambda_i u_i} u_i(1-u_i)f_t(u)]_0^{1-\sum_{j\neq i} u_j} \left(\prod_{j\neq i} e^{\lambda_j u_j} \mathbf{1}_{[0,1]}(u_j)du_j \right) -$$

$$\sum_{i=1}^{k} \int [e^{\lambda_i u_i}\partial_i \{u_i(1-u_i)f_t(u)\}]_0^{1-\sum_{j\neq i} u_j} \left(\prod_{j\neq i} e^{\lambda_j u_j} \mathbf{1}_{[0,1]}(u_j)du_j \right) -$$

$$\sum_{1\leq j\neq i\leq k} \int (\lambda_j u_j)[e^{\lambda_i u_i} u_i f_t(u)]_0^{1-\sum_{m\neq i} u_m} \left(\prod_{m\neq i} e^{\lambda_m u_m} \mathbf{1}_{[0,1]}(u_m)du_m \right) +$$

$$\sum_{1\leq j\neq i\leq k} \int [e^{\lambda_j u_j} u_j \partial_i \{u_i f_t(u)\}]_0^{1-\sum_{m\neq j} u_m} \left(\prod_{m\neq j} e^{\lambda_m u_m} \mathbf{1}_{[0,1]}(u_m)du_m \right).$$

By Leibniz rule, the third and the last terms split into

$$\sum_{i=1}^{k} \int [e^{\lambda_i u_i} u_i(1-u_i)\partial_i f_t(u)]_0^{1-\sum_{j\neq i} u_j} \left(\prod_{j\neq i} e^{\lambda_j u_j} \mathbf{1}_{[0,1]}(u_j)du_j \right) +$$

$$\sum_{i=1}^{k} \int [e^{\lambda_i u_i}(1-2u_i)f_t(u)]_0^{1-\sum_{j\neq i} u_j} \left(\prod_{j\neq i} e^{\lambda_j u_j} \mathbf{1}_{[0,1]}(u_j)du_j \right)$$

and

$$\sum_{1\leq j\neq i\leq k} \int u_i[e^{\lambda_j u_j} u_j \partial_i f_t(u)]_0^{1-\sum_{m\neq j} u_m} \left(\prod_{m\neq j} e^{\lambda_m u_m} \mathbf{1}_{[0,1]}(u_m)du_m \right) +$$

$$(k-1)\sum_{j=1}^{k} \int [e^{\lambda_j u_j} u_j f_t(u)]_0^{1-\sum_{m\neq j} u_m} \left(\prod_{m\neq j} e^{\lambda_m u_m} \mathbf{1}_{[0,1]}(u_m)du_m \right)$$

respectively. Thus there are no boundary terms at $u_i = 0, 1 \leq i \leq l$, while the remaining ones are given by

$$\sum_{i=1}^{k} \int \left[e^{\lambda_i u_i}(k+1-N)u_i f_t(u) \right]^{u_i = 1 - \sum_{j \neq i} u_j} \left(\prod_{j \neq i} e^{\lambda_j u_j} \mathbf{1}_{[0,1]}(u_j) du_j \right) +$$

$$\sum_{i=1}^{k} \int \left[e^{\lambda_i u_i} u_i \left\{ \lambda_i (1-u_i) - \sum_{j \neq i} \lambda_j u_j \right\} f_t(u) \right]^{u_i = 1 - \sum_{j \neq i} u_j} \left(\prod_{j \neq i} e^{\lambda_j u_j} \mathbf{1}_{[0,1]}(u_j) du_j \right) -$$

$$\sum_{i=1}^{k} \int \left[e^{\lambda_i u_i} u_i \left\{ (1-u_i) \partial_i f_t(u) - \sum_{j \neq i} u_j \partial_j f_t(u) \right\} \right]^{u_i = 1 - \sum_{j \neq i} u_j} \left(\prod_{j \neq i} e^{\lambda_j u_j} \mathbf{1}_{[0,1]}(u_j) du_j \right).$$

If $N \geq k+2$ and $f_t = g_t s_k$ vanishes on the hyperplane $\{u_1 + \cdots + u_k = 1\}$ and the boundary terms reduce to

$$\sum_{i=1}^{k} \int e^{\lambda_i \left(1 - \sum_{j \neq i} u_j \right)} \left(1 - \sum_{j \neq i} u_j \right) \left\{ \sum_{j \neq i} u_j [\partial_i f_t(u) - \partial_j f_t(u)] \right\}^{u_i = 1 - \sum_{j \neq i} u_j} \left(\prod_{j \neq i} e^{\lambda_j u_j} \mathbf{1}_{[0,1]}(u_j) du_j \right).$$

But since $\partial_i s_k = \partial_j s_k$ and since s_k vanishes on $\{u_1 + \cdots + u_k = 1\}$, then for any $1 \leq i \neq j \leq k$

$$\partial_i f_t(u) = \partial_j f_t(u), \quad u_1 + \cdots + u_k = 1$$

so that all boundary terms vanish. When $k = N - 1$ the boundary terms read

$$\sum_{i=1}^{k} \int e^{\lambda_i \left(1 - \sum_{j \neq i} u_j \right)} \left(1 - \sum_{j \neq i} u_j \right) \left\{ \sum_{j \neq i} (\lambda_i - \lambda_j) u_j \right\} [f_t(u)]^{u_i = 1 - \sum_{j \neq i} u_j} \left(\prod_{j \neq i} e^{\lambda_j u_j} \mathbf{1}_{[0,1]}(u_j) du_j \right) -$$

$$\sum_{i=1}^{k} \int e^{\lambda_i \left(1 - \sum_{j \neq i} u_j \right)} \left(1 - \sum_{j \neq i} u_j \right) \left\{ \sum_{j \neq i} u_j [\partial_i f_t(u) - \partial_j f_t(u)] \right\}^{u_i = 1 - \sum_{j \neq i} u_j} \left(\prod_{j \neq i} e^{\lambda_j u_j} \mathbf{1}_{[0,1]}(u_j) du_j \right).$$

References

[1] R. Aktas et Y. Xu – "Sobolev orthogonal polynomials on a simplex", *Int. Math. Res. Notice* **13** (2013), p. 3087–3131.

[2] G. E. Andrews, R. Askey et R. Roy – *Special functions*, Cambridge University Press, Cambridge, 1999.

[3] D. Bakry – "Remarques sur les semi-groupes de Jacobi", *Hommage à P. André Meyer et J. Neveu. Astérisque* **236** (1996), p. 23–39.

[4] A. Benabdallah – "Noyau de diffusion sur les espaces homogènes compacts", *Bull. Soc. Math. France* **101** (1973), p. 265–283.

[5] M. Berger, P. Gauduchon et E. Mazet – "Le spectre d'une variété Riemannienne", Lecture Notes in Mathematics, Springer-Verlag, 1971, p. vii+251 pp.

[6] C. Dunkl et Y. Xu – "Orthogonal Polynomials of Several Variables", Encyclopedia of Mathematics and Its Applications, Cambridge University Press, Cambridge, 2001.

[7] G. Gasper – "Banach algebras for Jacobi series and positivity of a kernel", *Ann. Math.* **95** (1972), p. 261–280.

[8] E. L. Grinberg – "Spherical harmonics and integral geometry on projective spaces", *Trans. Amer. Math. Soc.* **279** (1983), p. 187–203.

[9] F. Hiai et D. Petz – *The Semicircle Law, Free Random Variables and Entropy*, vol. 77, A. M. S., 2000.

[10] S. P. Karlin et G. McGregor – "Classical diffusion processes and total positivity", *J. Math. Anal. Appl.* **1** (1960), p. 163–183.

[11] T. Koornwinder – "The addition formula for Jacobi polynomials II. The Laplace type integral and the product formula", *Report TW 133/72. Mathematisch Centrum, Amsterdam* (1972).

[12] _____ , "The addition formula for Jacobi polynomials III. Completion of the proof", *Report TW 135/72. Mathematisch Centrum, Amsterdam* (1972).

[13] I. Nechita et C. Pellegrini – "Random pure quantum states via unitary Brownian motion", *Electron. Commun. Probab* **18** (2013), p. 1–13.

[14] N. Y. Vilenkin – *Fonctions spéciales et théorie de la représentation des groupes*, vol. 33, Monographies Universitaires de Mathématiques, Dunod, 1969.

[15] G. N. Watson – *A treatise on the theory of Bessel functions*, Cambridge University Press, 1995.

N. Demni

Nizar Demni
Institut de Recherche en Mathématiques
de Rennes
Université de Rennes 1
Campus de Beaulieu
35042 Rennes
FRANCE
nizar.demni@univ-rennes1.fr